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Measurement with the ATLAS detector of multi-particle azimuthal correlations in $p+{ m Pb}$ collisions at $\sqrt{s_{ m NN}}=5.02~{ m TeV}$

The ATLAS Collaboration

Abstract

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1 1. Introduction

Recent observations of ridge-like structures in 2 the two-particle correlation functions measured in 3 proton-lead (p+Pb) collisions at 5.02 TeV [1–3] 4 have led to differing theoretical explanations. These 5 structures have been attributed either to mecha-6 nisms that emphasize initial-state effects, such as the saturation of parton distributions in the Pb-8 nucleus [4-7], or to final-state effects, such as jet-9 medium interactions [8], interactions induced by 10 multiple partons [9–12], and collective anisotropic 11 flow [13–18]. 12

The collective flow of particles produced in nu-13 clear collisions, which manifests itself as a sig-14 nificant anisotropy in the plane perpendicular to 15 the beam direction, has been extensively studied 16 in heavy-ion experiments at the LHC [19–24] and 17 RHIC (for a review see Refs. [25, 26]). In p+Pb18 collisions the small size of the produced system 19 compared to the mean free path of the interacting 20 constituents might have been expected to generate 21 weaker collective flow, if any, compared to heavy-22 ion collisions. 23

However, two-particle correlation studies performed recently on data from p+Pb collisions at $\sqrt{s_{\rm NN}} = 5.02$ TeV revealed the presence of a

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"ridge", a structure extended in the relative pseudorapidity, $\Delta \eta$, while narrow in the relative azimuthal angle, $\Delta \phi$, on both the near-side ($\Delta \phi \sim 0$) [1] and away-side ($\Delta \phi \sim \pi$) [2, 3]. Furthermore, it was shown in Refs. [2, 3] that, after subtracting the component due to momentum conservation, the $\Delta \phi$ distribution in high-multiplicity interactions exhibits a predominantly $\cos(2\Delta \phi)$ shape, resembling the elliptic flow modulation of the $\Delta \phi$ distributions in Pb+Pb collisions.

The final-state anisotropy is usually characterized by the coefficients, v_n , of a Fourier decomposition of the event-by-event azimuthal angle distribution of produced particles [25, 27]:

$$v_n = \langle \cos n(\phi - \Psi_n) \rangle, \tag{1}$$

where ϕ is the azimuthal angle of the particle, Ψ_n is the event-plane angle for the *n*-th harmonic, and the outer brackets denote an average over charged particles in an event. In non-central heavy-ion collisions, the large and dominating v_2 coefficient is associated mainly with the elliptic shape of the nuclear overlap, and Ψ_2 defines the direction which nominally points in the direction of the classical impact parameter. In practice, initial-state fluctuations can blur the relationship between Ψ_2 and

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the impact parameter direction in nucleus-nucleus 52 collisions. In contrast, Ψ_2 in proton-nucleus would 53 be unrelated to the impact parameter and deter-54 mined by the initial-state fluctuations. In nucleus-55 nucleus collisions, the v_2 coefficient in central col-56 lisions and the other v_n coefficients in all collisions 57 are related to various geometric configurations aris-58 ing from fluctuations of the nucleon positions in the 59 overlap region [28]. 60

In this Letter, a direct measurement of the 61 second-order anisotropy parameter, v_2 , is presented 100 62 for p+Pb collisions at $\sqrt{s_{\rm NN}} = 5.02$ TeV. The cu-63 mulant method [29–32] is applied to derive v_2 using 102 64 two- and four-particle cumulants. The cumulant 65 method has been developed to characterize true 104 66 multi-particle correlations related to the collective 105 67 expansion of the system, while suppressing correla-68 tions from resonance decays, Bose-Einstein corre- 107 69 lations and jet production. Emphasis is placed on 108 70 the estimate of v_2 , $v_2\{4\}$, obtained from the four- 109 71 particle cumulants which are expected to be free 110 72 from the effects of short-range two-particle correla-111 73 tions, e.g. from resonance decays, unlike the two- 112 74 particle cumulants, used to estimate $v_2\{2\}$. 75

The measurements of multi-particle cumulants ¹¹⁴ 76 presented in this Letter should provide further con-¹¹⁵ 77 straints on the origin of long-range correlations ob-78 served in p+Pb collisions. 79

2. Event and track selections 80

The p+Pb data sample was collected during a 122 81 short run in September 2012, when the LHC deliv-123 82 ered p+Pb collisions at the nucleon–nucleon centre-124 83 of-mass energy $\sqrt{s_{\rm NN}} = 5.02$ TeV with the centre-84 of-mass frame shifted by -0.47 in rapidity relative ¹²⁶ 85 to the nominal ATLAS coordinate frame¹. 86 The measurements were performed using the ¹²⁸ 87

ATLAS detector [33]. The inner detector (ID) 88 was used for measuring trajectories and momenta 89

of charged particles for $|\eta| < 2.5$ with the silicon pixel detector and silicon microstrip detectors (SCT), and a transition radiation tracker, all placed in a 2 T axial magnetic field. For event triggering, two sets of Minimum Bias Trigger Scintillators (MBTS), located symmetrically in front of the endcap calorimeters, at $z = \pm 3.6$ m and covering the pseudorapidity range 2.1 < $|\eta| < 3.9$, were used. The trigger used to select minimumbias p+Pb collisions requires a signal in at least two MBTS counters. This trigger is fully efficient for events with more than four reconstructed tracks with $p_{\rm T} > 0.1$ GeV. The forward calorimeters (FCal), consisting of two symmetric systems with tungsten and copper absorbers and liquid argon as the active material, cover $3.1 < |\eta| < 4.9$ and are used to characterize the overall event activity.

The event selection follows the same requirements as used in the recent two-particle correlation analysis [3]. Events are required to have a reconstructed vertex with its z position within $\pm 150 \text{ mm}$ of the nominal interaction point. Beam–gas and photonuclear interactions are suppressed by requiring at least one hit in a MBTS counter on each side of the interaction point and at most a 10 ns difference between times measured on the two sides to eliminate through-going particles. To eliminate multiple p+Pb collisions (about 2% of collision events have more than one reconstructed vertex), the events with two reconstructed vertices that are separated in z by more than 15 mm are rejected. In addition, for the cumulant analysis presented here, it is required that the number of reconstructed tracks per event, passing the track selections as described below, is greater than three. With all the above selections, the analysed sample consists of about 1.9×10^6 events.

Charged particle tracks are reconstructed in the ID using the standard algorithm optimized for p+pminimum-bias measurements [34]. Tracks are required to have at least six hits in the SCT detector and at least one hit in the pixel detector. A hit in the first pixel layer is also required when the track crosses an active pixel module in that layer. Additional requirements are imposed on the transverse (d_0) and longitudinal $(z_0 \sin \theta)$ impact parameters measured with respect to the primary vertex. These are: $|d_0|$ and $|z_0 \sin \theta|$ must be smaller than 1.5 mm and must satisfy $|d_0/\sigma_{d_0}| < 3$ and $|z_0 \sin \theta/\sigma_z| < 3$, where σ_{d_0} and σ_z are uncertainties on the transverse and longitudinal impact parameters, respectively, as obtained from the covariance matrix of

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¹ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z-axis along the beam pipe. The x-axis points from the IP to the centre of the LHC ring, and the $y\text{-}\mathrm{axis}$ points upward. Cylindrical coordinates (r,ϕ) are used in the transverse plane, ϕ being the azimuthal angle around 137 the beam pipe. For the p+Pb collisions, the incident Pb beam travelled in the +z direction. The pseudorapidity is defined in laboratory coordinates in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$. Transverse momentum and energy are defined as $p_{\rm T} = p \sin \theta$ and $E_{\rm T} = E \sin \theta$, respectively.



Fig. 1: Upper plot: the ΣE_{T}^{Pb} distribution with the six activity intervals indicated. Lower plot: the distribution of $N_{\rm ch}^{\rm rec}$ for each activity interval. The leftmost distribution corresponds to the interval with the lowest ΣE_{T}^{Pb} , etc.

the track fit. The analysis is restricted to charged 142 175 particles with $0.3 < p_{\rm T} < 5.0$ GeV and $|\eta| < 2.5$. 143 The tracking efficiency is evaluated using 177 144 HIJING-generated [35] p+Pb events that are fully 178 145 simulated in the detector using GEANT4 [36, 37], 179 146 and processed through the same reconstruction 180 147 software as the data. The efficiency for charged 181 148 hadrons is found to depend only weakly on the 182 149 event multiplicity and on $p_{\rm T}$ for transverse mo- 183 150 menta above 0.5 GeV. An efficiency of about 82% 184 151 is observed at mid-rapidity, $|\eta| < 1$, decreasing to 185 152 about 68% at $|\eta| > 2$. For low- $p_{\rm T}$ tracks, between 186 153 0.3 GeV and 0.5 GeV, the efficiency ranges from $_{187}$ 154 74% at $\eta = 0$ to about 50% for $|\eta| > 2$. The number 155 of reconstructed charged particle tracks, not cor-189 156 rected for tracking efficiency, is denoted by $N_{\rm ch}^{\rm rec}$. 157

The analysis is performed in different intervals 191 158 of $\Sigma E_{\rm T}^{\rm Pb}$, the sum of transverse energy measured ¹⁹² 159 in the FCal with $3.1 < \eta < 4.9$ in the direction of 193 160 the Pb beam with no correction for the difference in 194 161

$\Sigma E_{\rm T}^{\rm Pb}$	$\langle \Sigma E_{\rm T}^{\rm Pb} \rangle$	range in	$\langle N_{\rm ch}^{\rm rec} \rangle$
range		fraction	(RMS)
[GeV]	[GeV]	of events [%]	
> 80	93.7	0-1.9	134(31)
55-80	64.8	1.9 - 9.1	102(26)
40-55	46.7	9.1 - 20.0	80 (23)
25-40	31.9	20.0 - 39.3	60(20)
10-25	16.9	39.3-70.4	37(17)
< 10	4.9	70.4–100	16(11)

Table 1: Characterization of activity intervals as selected by $\Sigma E_{\rm T}^{\rm Pb}$. In the last column, the mean and RMS of the number of reconstructed charged particles with $|\eta| < 2.5$ and $0.3 < p_{\rm T} < 5$ GeV, $N_{\rm ch}^{\rm rec}$, is given for each activity interval.

response to electrons and hadrons. The distribution of $\Sigma E_{\mathrm{T}}^{\mathrm{Pb}}$ for events passing all selection criteria is shown in Fig. 1. These events are divided into six $\Sigma E_{\mathrm{T}}^{\mathrm{Pb}}$ intervals to study the variation of v_2 with overall event activity, as indicated in Fig. 1 and shown in Table 1. Event "activity" is characterized by $\Sigma E_{\rm T}^{\rm Pb}$: the most active events are those with the largest $\Sigma E_{\rm T}^{\rm Pb}$. The distribution of $N_{\rm ch}^{\rm rec}$ for each activity interval is shown in the lower plot of Fig. 1.

3. Data analysis

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The cumulant method involves the calculation of 2k-particle azimuthal correlations, $corr_n\{2k\}$, and cumulants, $c_n\{2k\}$, where k = 1, 2 for the analysis presented in this paper. The two- and four-particle correlations are defined as $corr_n\{2\} = \langle e^{in(\phi_1 - \phi_2)} \rangle$ and $corr_n\{4\} = \langle e^{in(\phi_1 + \phi_2 - \phi_3 - \phi_4)} \rangle$, respectively, where the angle brackets denote the average in a single event over all pairs and all combinations of four particles. After averaging over events, the two-particle cumulant is obtained as $c_n\{2\} = \langle corr_n\{2\} \rangle$, and the four-particle cumulant $c_n\{4\} = \langle corr_n\{4\} \rangle - 2 \cdot \langle corr_n\{2\} \rangle^2$. Thus the effect of two-particle correlations is explicitly removed in the expression for $c_n\{4\}$. Further details are given in Refs. [29, 30, 32].

Direct calculation of multi-particle correlations requires multiple passes over the particles in an event, and requires extensive computing time in high-multiplicity events. To mitigate this, it has been proposed in Ref. [32] to express multi-particle correlations in terms of the moments of the flow vector Q_n , defined as $Q_n = \sum_i e^{in\phi_i}$, where the index *n* denotes the flow harmonic and the sum runs

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over all particles in an event. This analysis is re-195

stricted to the second harmonic coefficient, n = 2. 196

The method based on the flow-vector moments en-197 ables the calculation of multi-particle cumulants in 198 a single pass over the full set of particles in each 199 event. 200

244 The cumulant method involves two main steps 201 245 [29, 30]. In the first step, the so-called "refer-202 246 ence" flow harmonic coefficients are calculated us-203 ing multi-particle cumulants for particles selected 204 inclusively from a broad range in $p_{\rm T}$ and η as: 205

$$v_2^{\text{ref}}\{2\} = \sqrt{c_2\{2\}},$$
 (2)

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$$v_2^{\text{ref}}\{4\} = \sqrt[4]{-c_2\{4\}},\tag{3}$$

where $v_2^{\text{ref}}\{2\}$ $(v_2^{\text{ref}}\{4\})$ denotes the reference es-206 timate of the second-order anisotropy parameter 207 obtained using two-particle, $c_2\{2\}$ (four-particle, 208 $c_2\{4\}$) cumulants. 209

The flow-vector method is easiest to apply when 210 the detector acceptance is azimuthally uniform [32]. 211 A correction for any azimuthal non-uniformity in 212 the reconstruction of charged particle tracks is ob-213 tained from the data [25], based on an $\eta - \phi$ map 214 of all reconstructed tracks. For each small ($\delta \eta =$ 215 $0.1, \delta \phi = 2\pi/64$) bin (labelled *i*), a weight is cal-216 culated as $w_i(\eta, \phi) = \langle N(\delta\eta) \rangle / N_i(\delta\eta, \delta\phi)$, where 217 $\langle N(\delta \eta) \rangle$ is the event-averaged number of tracks 218 in the $\delta\eta$ slice to which this bin belongs, while 219 $N_i(\delta\eta, \delta\phi)$ is the number of tracks in an event 220 within this bin. Using this weight forces the az-221 imuthal angle distribution of reference particles to 222 be uniform in ϕ , but it does not change the η 223 distribution of reconstructed tracks. A weighted 224 Q-vector is evaluated as $Q_n = \sum_i w_i e^{in\phi_i}$ [32]. 225 From Eqs. (2) and (3) it is clear that the cumu-226 lant method can be used to estimate v_2 only when 227 $c_2\{4\}$ is negative and $c_2\{2\}$ positive. 228

In the second step, the harmonic coefficients are 229 determined as functions of $p_{\rm T}$ and η , in bins in each 230 variable (10 bins of equal width are used in η and 22 231 bins of varied width in $p_{\rm T}$). These differential flow 232 harmonics are calculated for "particles of interest" 233 which fall into these small bins. First, the differen-234 tial cumulants, $d_2\{2\}$ and $d_2\{4\}$, are obtained by 235 correlating every particle of interest with one and 236 three reference particles respectively. The differen-237 tial second harmonic, $v_2\{2k\}(p_T, \eta)$, where k = 1, 2, 238 is then calculated with respect to the reference flow 239 as derived in Refs. [29, 30]: 240

$$v_2\{2\}(p_{\rm T},\eta) = \frac{d_2\{2\}}{\sqrt{c_2\{2\}}},$$
 (4) ²⁵⁸
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$$v_{2}\{4\}(p_{\mathrm{T}},\eta) = \frac{-d_{2}\{4\}}{\sqrt[3]{4}-c_{2}\{4\}}.$$
(5)

The differential v_2 harmonic is then integrated over wider phase-space bins, with each small bin weighted by the appropriate charged particle multiplicity. This is obtained from the reconstructed multiplicity by applying η - and $p_{\rm T}$ -dependent efficiency factors, determined from Monte Carlo (MC) simulation as discussed in the previous section. Due to the small number of events in the data sample, the final results are integrated over the full acceptance in η .



Fig. 2: The two-particle (upper plot) and four-particle (lower plot) cumulants calculated using the reference flow particles as a function of $\Sigma E_{\rm T}^{\rm Pb}$ for data (circles), the fully simulated HIJING events (open squares) and the large generator-level HIJING sample (filled squares). For clarity, the points for the fully simulated (generated) HIJING events are slightly shifted to the left (right).

Fig. 2 shows the two- and four-particle cumulants, averaged over events in each event-activity class defined in Table 1, as a function of $\Sigma E_{\rm T}^{\rm Pb}$. It is observed that four-particle cumulants are negative only in a certain range of event activity. This restricts subsequent analysis to events with $\Sigma E_{\rm T}^{\rm Pb} > 25$ GeV, for which the four-particle cumulant in data is found to be less than zero by at least two standard deviations (statistical errors only). It

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was also checked that for these events $c_2\{4\}$ is un-260 changed within errors for any high-multiplicity se- 313 261 lection. For example, defining N_{20} as the value of $_{314}$ 262 $N_{
m ch}^{
m rec}$ such that 20% of events have $N_{
m ch}^{
m rec}$ < $N_{
m 20}$ $_{
m 315}$ 263 (i.e. N_{20} is the 20th percentile), then selecting $_{316}$ 264 $N_{\rm ch}^{\rm rec} > N_{20}$ leaves $c_2\{4\}$ unchanged within errors. 317 265 And for $\Sigma E_{\rm T}^{\rm Pb} > 25$ GeV this holds for any per- 318 266 centile selection. 267 319

Fig. 2 also shows the cumulants calculated for 320 268 50 million HIJING-generated events, using the true 321 269 particle information only, as well as for one million 322 270 fully simulated and reconstructed HIJING events, 323 271 using the same methods as used for the data. The $_{\rm 324}$ 272 $\Sigma E_{\rm T}^{\rm Pb}$ obtained from the HIJING sample is rescaled 325 273 to match that measured in the data. It should be 326 274 noted that the HIJING Monte Carlo model does 327 275 not contain any collective flow, and the only corre- 328 276 lations are those due to resonance decays, jet pro-329 277 duction and momentum conservation. The values 330 278 of $c_2\{2\}$ for HIJING events are smaller than the ${}_{331}$ 279 values obtained from the data, and there is no sig- 332 280 nificant difference between the HIJING results ob-333 281 tained at the generator ("truth") level and at the 334 282 reconstruction level. For $c_2{4}$, the HIJING events 283 at $\Sigma E_{\rm T}^{\rm Pb} \sim 20 \,\,{\rm GeV}$ show a negative value compa-28 rable to the values seen in the data, indicating that 285 correlations from jets or momentum conservation 336 286 contribute significantly to $v_2{4}$ in events of low 337 287 multiplicity. For $\Sigma E_{\rm T}^{\rm Pb} > 25$ GeV the generator-level HIJING sample's values for $c_2\{4\}$ are also neg-288 289 ative, but the magnitude is much smaller than in 340 290 the data or in HIJING events with smaller ΣE_{T}^{Pb} . ³⁴¹ 291 The size of the fully simulated HIJING event sam- 342 292 ple is too small to draw a definite conclusion about 343 293 the sign or magnitude of $c_2\{4\}$. 294

The systematic uncertainties on $v_2\{2\}$ and $v_2\{4\}$ 345 295 as a function of $p_{\rm T}$ and $\Sigma E_{\rm T}^{\rm Pb}$ have been evaluated $_{346}$ 296 by varying several aspects of the analysis proce- 347 297 dure. Azimuthal-angle sine terms in the Fourier 348 298 expansion should be zero, but a non-zero contribu- 349 299 tion can arise due to detector biases. It was found 350 300 that the magnitude of the sine terms relative to 351 301 the cosine terms is negligible (below 1%) for v_2 {2} 352 302 measured as a function of $p_{\rm T}$, as well as for the $_{353}$ 303 $p_{\rm T}$ -integrated $v_2\{2\}$ and $v_2\{4\}$. In the case of the ${}_{354}$ 304 measurement of the $p_{\rm T}$ -dependent $v_2\{4\}$, the sys-355 305 tematic uncertainty attributed to the residual sine 356 306 terms varies between 6% and 14% in the different $_{357}$ 307 $\Sigma E_{\mathrm{T}}^{\mathrm{Pb}}$ intervals. Uncertainties related to the track-308 ing are obtained from the differences between the 359 309 main results and those using tracking requirements 360 310 modified to be either more or less restrictive. They 361 311

are found to be negligible (below 0.2%) for $v_2\{2\}$. For the $p_{\rm T}$ -dependent $v_2\{4\}$ they give a contribution of less than 6% to the systematic uncertainty, and less than 1% for the $p_{\rm T}$ -integrated v_2 {4}. In addition to varying the track quality requirements, an uncertainty on the $p_{\rm T}$ dependence of the efficiency corrections is also taken into account, and found to be below 1% for the v_2 {2} and v_2 {4} measurements. The correction of the azimuthal-angle uniformity is checked by comparing the results to those obtained with all weights, w_i , set equal to one. This change leads to small relative differences, below 1%, for the $v_2\{2\}$ measured as a function of p_T , as well as for the $p_{\rm T}$ -integrated $v_2\{2\}$ and $v_2\{4\}$. Up to 4% differences are observed in the $p_{\rm T}$ -dependent $v_2\{4\}$. All individual contributions to the systematic uncertainty are added in quadrature and quoted as the total systematic uncertainty. The total systematic uncertainties are below 1% for the $v_2\{2\}$ measurement. The $v_2{4}$ measurement precision is limited by large statistical errors, whereas the systematic uncertainties stay below 15% for v_2 {4}(p_T) and below 2% for the $p_{\rm T}$ -integrated v_2 {4}.

4. Results

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Fig. 3 shows the transverse momentum dependence of $v_2\{2\}$ and $v_2\{4\}$ in four different classes of the event activity, selected according to $\Sigma E_{\rm T}^{\rm Pb}$. A significant second-order harmonic is observed. $v_2\{4\}$ is systematically smaller than $v_2\{2\}$, consistent with the suppression of non-flow effects in v_2 {4}. This difference is most pronounced at high $p_{\rm T}$ and in collisions with low $\Sigma E_{\rm T}^{\rm Pb}$ where jet-like correlations not diluted by the underlying event can contribute significantly. Thus, v_2 {4} appears to provide a more reliable estimate of the second-order anisotropy parameter of collective flow. As a function of transverse momentum the second-order harmonic, v_2 {4}, increases with p_T up to $p_T \approx 2$ GeV. Large statistical errors preclude a definite conclusion about the $p_{\rm T}$ dependence of $v_2\{4\}$ at higher transverse momenta.

The shape and magnitude of the $p_{\rm T}$ -dependence of $v_2\{4\}$ is found to be similar to that observed in p+Pb collisions using two-particle correlations [2, 3]. The second-order harmonic, v_2 , can be extracted from two-particle azimuthal correlations using charged particle pairs with a large pseudorapidity gap to suppress the short-range correlations on the near-side $(\Delta \phi \sim 0)$ [3, 22]. However, the twoparticle correlation measured this way may still be



Fig. 3: The second-order harmonic calculated with the two-particle (circles) and four-particle (stars) cumulants as a function of transverse momentum in four different activity intervals. Bars denote statistical errors; systematic uncertainties are shown as shaded bands. The v_2 derived from the Fourier decomposition of two-particle correlations [3] is shown by squares.

affected by the dijet correlations on the away-side 362 $(\Delta \phi \sim \pi)$, which can span a large range in $\Delta \eta$. 363 In Ref. [3], the away-side non-flow correlation is 364 estimated using the yield measured in the lowest 365 $\Sigma E_{\rm T}^{\rm Pb}$ collisions and is then subtracted from the 366 higher $\Sigma E_{\rm T}^{\rm Pb}$ collisions. The result of that study, 367 $v_2\{2PC\}$, is shown in Fig. 3 for the four activ-368 ity intervals with largest $\Sigma E_{\rm T}^{\rm Pb}$, and compared to 369 v_2 {4}. Good agreement is observed between v_2 {4} 370 and $v_2\{2PC\}$ for collisions with $\Sigma E_T^{Pb} > 55$ GeV. 371 For $\Sigma E_{\rm T}^{\rm Pb}$ < 55 GeV, the disagreement could be 372 due either to the subtraction procedure used to ob-373 tain $v_2\{2PC\}$ or to non-flow effects in $v_2\{4\}$, or to 374 a combination. 375

The dependence on the collision activity of the 376 second-order harmonic, integrated over $0.3 < p_{\rm T} <$ 377 5 GeV, is shown in Fig. 4. The large magni-378 tude of $v_2\{2\}$ compared to $v_2\{4\}$ suggests a sub-379 stantial contamination from non-flow correlations. 380 The value of $v_2{4}$ is approximately 0.06, with lit-381 tle dependence on the overall event activity for 382 $\Sigma E_{\rm T}^{\rm Pb}$ >25 GeV. The extracted values of v_2 {4} 383 are also compared to the $v_2\{2PC\}$ values obtained 384 from two-particle correlations. Good agreement is 395 385 observed at large $\Sigma E_{\rm T}^{\rm Pb}$, while at lower $\Sigma E_{\rm T}^{\rm Pb}$ the ³⁹⁶ $v_2\{2PC\}$ is smaller than $v_2\{4\}$, which may be due ³⁹⁷ 386 387 to different sensitivity of the two methods to non- 398 388 flow contributions that become more important in 399 389 low $\Sigma E_{\rm T}^{\rm Pb}$ collisions. Although v_2 {4} is constructed 400 390 to suppress local two-particle correlations, it may 401 391 still include true multi-particle correlations from $_{402}$ 392 jets, which should account for a larger fraction of 403 393 the correlated particle production in the events with 404 394



Fig. 4: The second-order harmonic, v_2 , integrated over p_T and η , calculated with two- and four-particle cumulants (circles and stars, respectively), as a function of ΣE_T^{Pb} . Systematic uncertainties are shown as shaded bands. Also shown is $v_2\{2PC\}$ (squares) and predictions from the hydrodynamic model [18] (triangles) for the same selection of charged particles as in the data.

the lowest $\Sigma E_{\mathrm{T}}^{\mathrm{Pb}}$. If the HIJING results, shown in Fig. 2, were used to correct the measured cumulants for this non-flow contribution, the extracted $v_2\{4\}$ would be decreased by at most 10% for $v_2\{4\}$ shown in Fig. 4. However, this correction is not applied to the final results.

It is notable that the trend of the $p_{\rm T}$ dependence of both $v_2\{4\}$ and $v_2\{2PC\}$ in $p+{\rm Pb}$ collisions resembles that observed for v_2 measured with the event-plane method in Pb+Pb collisions

at $\sqrt{s_{\rm NN}} = 2.76$ TeV [21, 22], although with a mag-405 nitude between that observed in the most central 449 406 and peripheral Pb+Pb collisions. While the trend 450 407 is found to be nearly independent of the Pb+Pb 451 408 collision geometry, the magnitude in Pb+Pb events 452 409 depends on the initial shape of the colliding sys-410 tem, and has been modelled for $p_{\rm T} < 2$ GeV using 411 453 viscous hydrodynamics [39–41]. 412

Harmonic flow coefficients in p+Pb collisions at 413 $\sqrt{s_{\rm NN}} = 5.02$ TeV have also been predicted using 414 viscous hydrodynamics, with similar initial condi-415 tions as the Pb+Pb calculations [18]. The pre-416 dicted magnitude of the second-order $\rm harmonic^2$ 417 458 is compared to the measured $v_2{4}$ and $v_2{2PC}$ 418 459 in Fig. 4. It can be seen that the hydrodynamic 419 predictions agree with our measurements over the 420 $\Sigma E_{\rm T}^{\rm Pb}$ range where the model predictions are avail-421 able. 422

5. Conclusions 423

ATLAS has measured the second harmonic coef-424 ficient in p+Pb collisions at $\sqrt{s_{\rm NN}} = 5.02$ TeV us-425 ing two- and four-particle cumulants. A significant 426 magnitude of v_2 is observed using both two- and 427 four-particle cumulants, although $v_2\{2\}$ is consis-428 tently larger than $v_2\{4\}$, indicating a sizeable con-429 tribution of non-flow correlations to $v_2\{2\}$. The 430 transverse momentum dependence of $v_2{4}$ shows 431 a behaviour similar to that measured in Pb+Pb 432 collisions at $\sqrt{s_{\rm NN}} = 2.76$ TeV. The magnitude of 433 v_2 {4} increases with p_T up to about 2–3 GeV. As 434 a function of the collision activity, v_2 {4} remains 435 constant, at the level of about 0.06, for the colli-436 sions with $\Sigma E_{\rm T}^{\rm Pb} > 25$ GeV, which corresponds to 437 about 40% of the data. The measured $v_2{4}$ is 438 found to be consistent with the second harmonic co-439 efficient extracted by the Fourier decomposition of 440 the long-range two-particle correlation function for 441 collisions with $\Sigma E_{\rm T}^{\rm Pb} > 55$ GeV. Good agreement is 442 also found with the predictions of a hydrodynamic 443 calculation for p+Pb collisions. 444

Extending previous results based only on two-445 particle correlations, the multi-particle cumulant 446 results presented here provide additional evidence 447

for the importance of final-state effects in the highest multiplicity p+Pb reactions. Final-state effects may lead to collective flow similar to that observed in the hot, dense system created in high-energy heavy-ion collisions.

6. Acknowledgements

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